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# Modern approaches to quantum gravity

## Homework 6

Fall 2025

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### 1. A stack of strings

- (a) State the Polyakov and Nambu-Goto actions for a string coupled to a  $D$ -dimensional background metric  $G_{\mu\nu}$ ,  $S_{\text{Poly}}[G]$  and  $S_{NG}[G]$ . Assume that no background  $B$ -field or dilaton  $\Phi$  are present.
- (b) The string frame metric produced by  $N$  infinite static strings, i.e. strings, lying in the  $(X^0, X^1) \equiv (t, x)$  plane is

$$ds^2 = f(r)^{-1}(-dt^2 + dx^2) + d\vec{X} \cdot d\vec{X} \quad (1)$$

where  $\vec{X} = (X^2, \dots, X^{D-1})$  labels the directions transverse to the strings and

$$f(r) = 1 + g_s^2 N \left(\frac{l_s}{r}\right)^{D-4}, \quad r^2 \equiv \vec{X} \cdot \vec{X}. \quad (2)$$

Consider one further infinite probe string in this background, lying parallel to the others, separated by a distance  $r$ .

Write down the Nambu-Goto action describing the motion of this string. Show that in static gauge  $t = \tau$  and  $x = \sigma$ , the low-energy excitations of the string are governed by the effective action

$$L \approx \frac{1}{2\pi\alpha'} \int dt dx \left[ -f(r)^{-1} + \frac{1}{2} \left( \frac{d\vec{X}}{dt} \cdot \frac{d\vec{X}}{dt} - \frac{d\vec{X}}{dx} \cdot \frac{d\vec{X}}{dx} \right) + \dots \right] \quad (3)$$

Find the equations of motion of the string and interpret the result.

If the string is also coupled to an external  $B$ -field, the worldsheet action is modified by adding

$$S = S[G] + \frac{1}{4\pi\alpha'} \int d^2\sigma \sqrt{g} B_{\mu\nu} \partial_\alpha X^\mu \partial_\beta X^\nu \epsilon^{\alpha\beta} \quad (4)$$

where we introduced the antisymmetric tensor density  $\epsilon^{\alpha\beta}$ , normalised so that  $\sqrt{g}\epsilon^{01} = 1$ .

- (c) In the setup from the previous point, the  $N$  strings produce a background  $B$ -field with a single independent component

$$B_{01} = f(r)^{-1} - 1 \quad (5)$$

Show that the probe string feels no static force.

## 2. The different regimes of the Dp-brane gravity description

Before tackling this exercise, let us give a reminder on the different parameters of string theory. The Regge slope  $\alpha' = 1/M_s^2 = \ell_s^2$  sets the scale at which string states become important. The string coupling is  $e^\phi = g_s$ , and it relates to the Planck mass as  $g_s \sim M_s^4/M_p^4$ , thus  $g_s$  controls the size of perturbative quantum gravity corrections.

Let us consider the 10-dimensional supergravity solution describing  $N$  coincident extremal Dp-branes in the string frame

$$ds^2 = f_p^{-1/2}(-dt^2 + dx_1^2 + \dots dx_p^2) + f_p^{1/2}(dx_{p+1}^2 + \dots + dx_9^2) \quad (6)$$

$$e^{-2(\phi-\phi_\infty)} = f_p^{(p-3)/2} \quad (7)$$

where

$$\alpha'^2 f_p = \alpha'^2 + \frac{d_p g_{\text{YM}}^2 N}{U^{7-p}} \quad d_p = 2^{7-2p} \pi^{\frac{9-3p}{2}} \Gamma\left(\frac{7-p}{2}\right) \quad (8)$$

with  $U = r/\alpha'$ ,  $r$  being the radial variable for the transverse directions  $x_{p+1}, \dots, x_9$ . We will consider the field theory limit

$$U = \text{fixed} \quad g_{\text{YM}}^2 = (2\pi)^{p-2} g_s \alpha'^{(p-3)/2} = \text{fixed} \quad \alpha' \rightarrow 0 \quad (9)$$

The parameter  $U$  is interpreted, in the gauge theory side, as the energy scale which we want to keep finite.

- (a) What does the limit (9) imply on  $g_s$  when  $p > 3$ ? What is special about the case  $p = 3$ ?
- (b) In the limit (9), show that the metric and the dilaton  $\phi$  reduce to

$$ds^2 = \alpha' \left( \frac{U^{(7-p)/2}}{g_{\text{YM}} \sqrt{d_p N}} dx_{\parallel}^2 + \frac{g_{\text{YM}} \sqrt{d_p N}}{U^{(7-p)/2}} dU^2 + g_{\text{YM}} \sqrt{d_p N} U^{(p-3)/2} d\Omega_{8-p}^2 \right) \quad (10)$$

$$e^\phi = (2\pi)^{(2-p)} g_{\text{YM}}^2 \left( \frac{g_{\text{YM}}^2 d_p N}{U^{7-p}} \right)^{\frac{3-p}{4}} \quad (11)$$

- (c) Show that the curvature (Ricci scalar) associated to this metric is given by

$$\alpha' R \sim \frac{1}{g_{\text{eff}}} \quad (12)$$

where we defined the effective coupling in the field theory side

$$g_{\text{eff}}^2 = \frac{g_{\text{YM}}^2 N}{U^{3-p}} \quad (13)$$

- (d) The supergravity regime is trustworthy whenever the coupling is small  $g_s \sim e^\phi \ll 1$  and the curvature is small compared to the string scale  $\alpha' R \ll 1$ . Show that satisfying both conditions implies

$$1 \ll g_{\text{eff}}^2 \ll N^{\frac{4}{7-p}} \quad (14)$$

- (e) Rewrite the above condition in terms of  $U$ , and distinguish physically each region. In particular, what happens to the dilaton at large radius and at small radius for different values of  $p$ ?